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BUILDING A SCIENTIFIC FOUNDATION FOR SOUND ENVIRONMENTAL DECISIONS

In Pursuit of the Elusive Bound Site Activity Coefficient

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Introduction

Diffuse Layer Models (DLMs) are widely employed to describe the acidity behavior of bound reactive sites on metal oxides, NOM, polyelectrolytes and latex beads:



$$K_{a1} = \frac{[>SO^-]a_{H^+}EXP(-e\psi/kT)}{[>SOH]}$$

Historically, these mixed concentration/activity expressions assume bound site activity coefficients: 1) cancel out in the quotient, 2) equal 1 (i.e., ideal behavior), or 3) are represented in the exponential term.

Following Chan et al. (1975), the electrochemical potentials of these species are given by:

$$u_{\text{H}^+} = u_{\text{H}^+}^{\circ} + kT \ln([\text{H}^+]) + kT \ln(\gamma_{\text{H}^+}) - e\psi$$

$$u_{>\text{SOH}} = u_{<\text{SOH}}^{\circ} + kT \ln([>\text{SOH}]) + kT \ln(\gamma_{>\text{SOH}})$$

$$u_{>\text{SO}^-} = u_{>\text{SO}^-}^{\circ} + kT \ln([>\text{SO}^-]) + kT \ln(\gamma_{>\text{SO}^-}) - e\psi$$

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Due to moving a mobile ion from bulk solution where $\Psi=0$ to the interface where $\Psi \neq 0$.

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Due to the creation of a charged site in a region of constant potential Ψ .

Problem: If one creates a mass action expression from these electrochemical potentials (assuming that charging energies are the only significant non-ideal excess free energy term), one arrives at:

$$K_{a1} = \frac{[>SO^-]a_{H^+}EXP(-2e\psi/kT)}{[>SOH]}$$

where 1 e Ψ is due to the hydronium ion and 1 e Ψ is due to the creation of a charged site in a region of constant potential Ψ .

Incorporation of counterion condensation into the electrochemical potentials:

$$u_{\text{H}^+} = u_{\text{H}^+}^{\circ} + kT \ln([\text{H}^+]) + kT \ln(\gamma_{\text{H}^+}) - e\Psi$$

$$u_{>\text{SOH}} = u_{<\text{SOH}}^{\circ} + kT \ln([>\text{SOH}]) + kT \ln(\gamma_{>\text{SOH}})$$

$$u_{>\text{SO}^-} = u_{>\text{SO}^-}^{\circ} + kT \ln([>\text{SO}^-]) + kT \ln(\gamma_{>\text{SO}^-}) - (1 - \tau) e\Psi$$

Loux (1999, 2000, 2006) modified the bound site charging energy term to include counterion condensation (τ ; where $0 \leq \tau \leq 1$)

Assuming that charging energies are the only significant non-ideal excess free energy term, at high charge densities and constant τ the new mass action expression is:

$$K_{a1} = \frac{[>SO^-]a_{H^+}EXP(-(2-\tau)e\psi/kT)}{[>SOH]}$$

where $(1-\tau)e\psi$ is the new charging energy term. Note that as τ approaches 1, this expression collapses into the historical approach.

Activity coefficients for ions in bulk solution are obtained from charging energies in the Debye-Huckel Limiting Law (DHLL) for spherical ions. The DHLL is derived under the following assumptions:

1) $e\psi \ll kT$

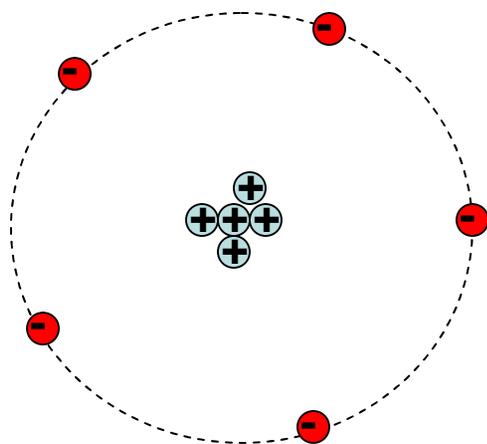
2) $1/\kappa \gg r_{ion}$

3) the charge in the diffuse layer can be represented as a hollow charged sphere $1/\kappa$ distant from the surface of the ion.

DHLL

(Hunter, 1987)

$$\left| -1/\kappa \right|$$



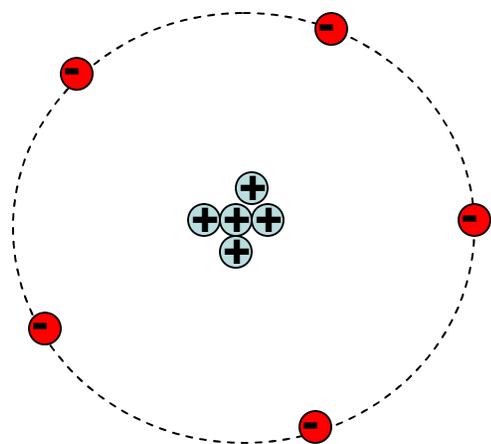
$$e\psi_d \ll kT$$

$$\psi_d = \frac{Q_p}{4\pi\epsilon r_{pcl}} - \frac{Q_p}{4\pi\epsilon(r_{pcl} + 1/\kappa)}$$

DHLL

(Hunter, 1987)

$$\left| -\frac{1}{\kappa} \right|$$

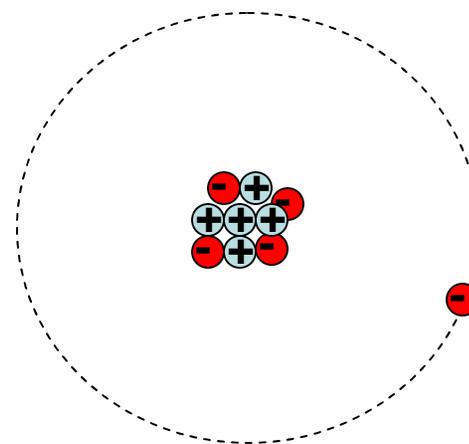


$$e\Psi_d \ll kT$$

$$\Psi_d = \frac{Q_p}{4\pi\epsilon r_{pcl}} - \frac{Q_p}{4\pi\epsilon(r_{pcl} + 1/\kappa)}$$

Present work

$$\left| -\frac{1}{\kappa} \right|$$



$$e\Psi_d \gg kT$$

$$\Psi_d = \frac{Q_p}{4\pi\epsilon r_{pcl}} - \frac{\tau Q_p}{4\pi\epsilon r_{pcl}}$$

Estimation of the charging energy terms:

DHLL

$$\Delta G = \int_0^{ze} [\kappa Q_d / 4\pi\epsilon] dQ$$

$$= (ze)^2 \kappa / 8\pi\epsilon$$

Present work

$$\Delta G = \int_0^{Q \pm e} [(1-\tau)Q / 4\pi\epsilon r_{pcl}] dQ$$

$$= (1-\tau)(\pm 2eQ + e^2) / 8\pi\epsilon r_{pcl}$$

$$= \pm(1-\tau)e\Psi$$

Assumptions:

$$e\Psi \ll kT$$

$$1/\kappa \gg r_{ion}$$

Q >> e and constant τ , most charge neutralization from CC



Estimation of the activity coefficients:

DHLL

$$\ln(\gamma) = -\Delta G/kT$$

$$= -(ze)^2\kappa/8\pi\epsilon kT$$

Present work

$$\ln(\gamma) = -\Delta G/kT$$

$$= -/+ (1-\tau)e\Psi/kT$$

Comparison between the DHLL and the present approach.

In applications of the the DHLL, $\tau \neq 0$; e.g., for a 1:1 ion pair, $\tau = 1$, $\Delta G = 0$ and $\gamma = 1$. Therefore, a more accurate DHLL potential expression (with no change in computational accuracy) is given by:

$$\Psi_d = \frac{(1-\tau)Q_p}{4\pi\epsilon r_{pcl}} - \frac{(1-\tau)Q_p}{4\pi\epsilon(r_{pcl} + 1/\kappa)}$$

Or the present approach is a truncation where:

$$|- \tau Q_p / 4\pi\epsilon r_{pcl}| \gg |-(1-\tau)\kappa Q_p / 4\pi\epsilon|$$

(i.e., CC dominates charge neutralization).



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Methods

Surface potential estimates for spherical particles were obtained using the approximate analytical solution to the P-B equation published by Oshima et al. (1982). $T = 298$ K.

Following Loux (2006): $T_{PB} = (\Psi_{Coul.} - \Psi_{PB})/\Psi_{Coul.}$

Because τ monotonically approaches a maximum value with increased charge density, maximum values of τ (or τ_{max}) were estimated using a Michaelis-Menten type plot of σ/τ vs. σ where the slope equals $1/\tau_{max}$.

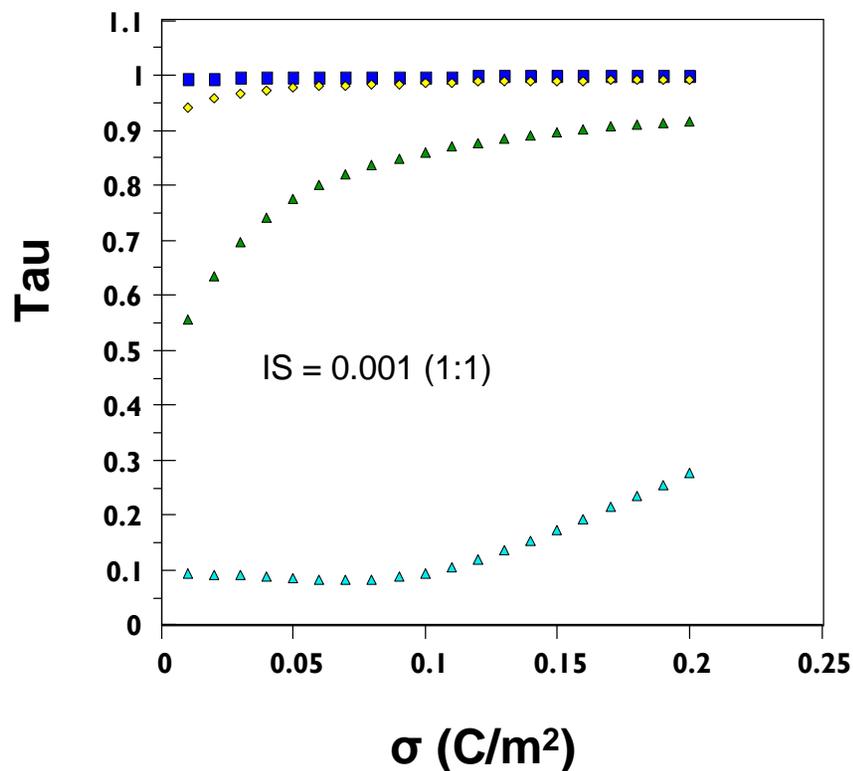


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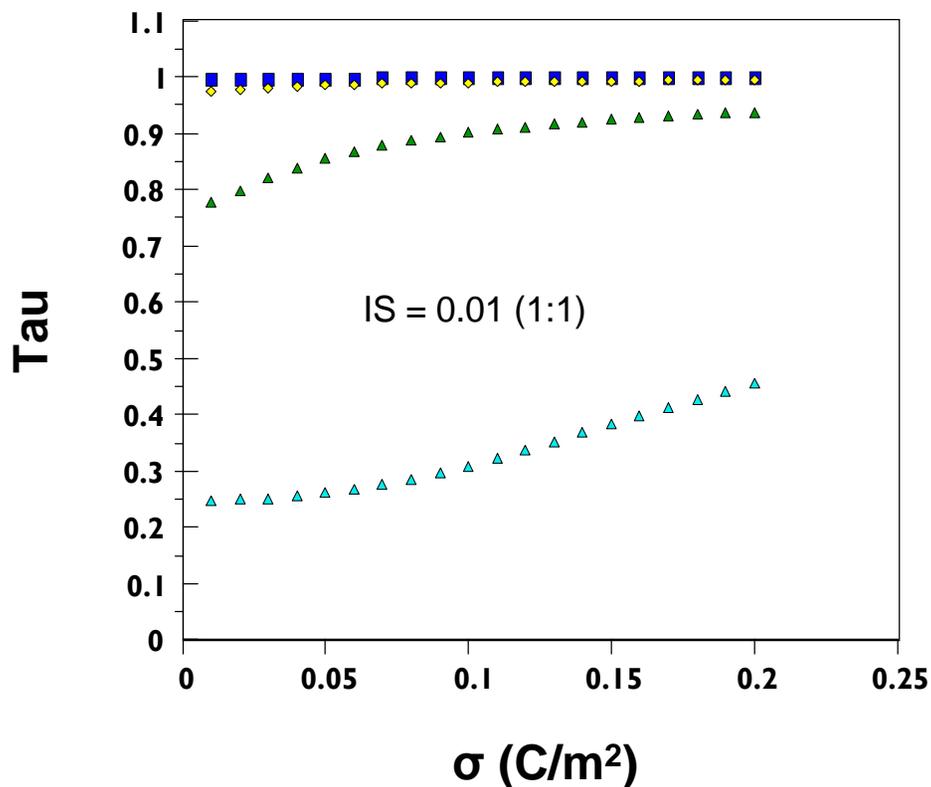
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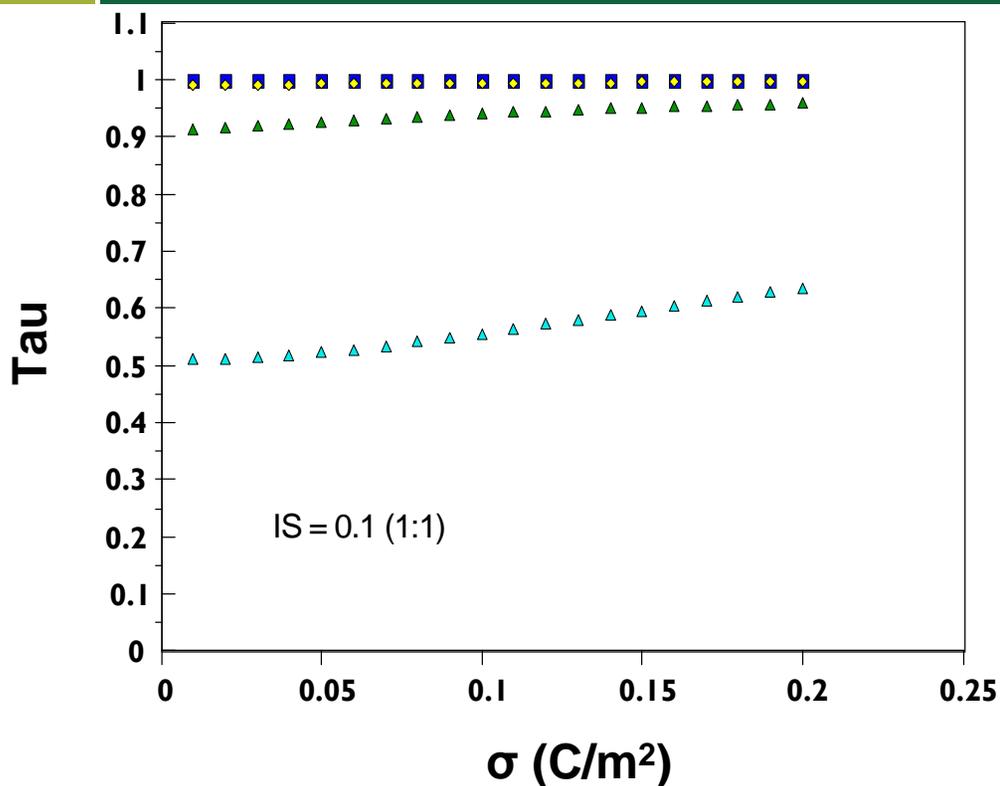
Results



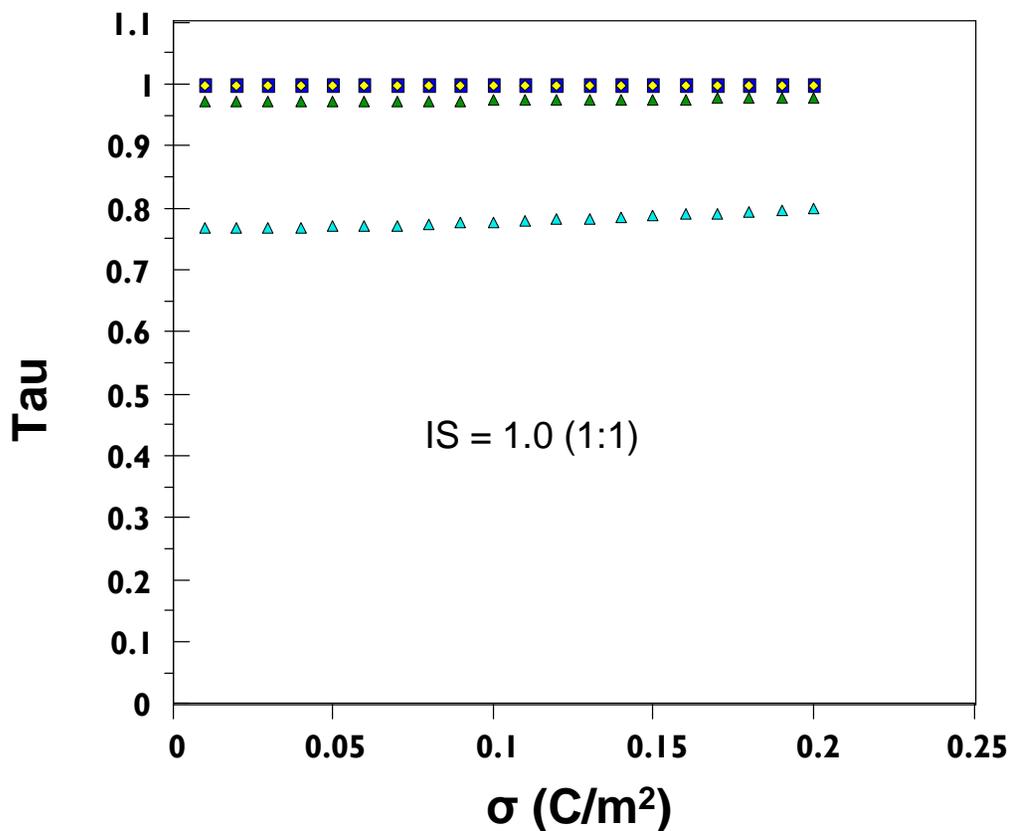
“ τ ” as a f(r_{pcl} and σ): $r_{pcl} = 1.0 \mu\text{m}$ (blue boxes), $0.1 \mu\text{m}$ (yellow diamonds), $0.01 \mu\text{m}$ (green triangles), and $0.001 \mu\text{m}$ (aquamarine triangles).



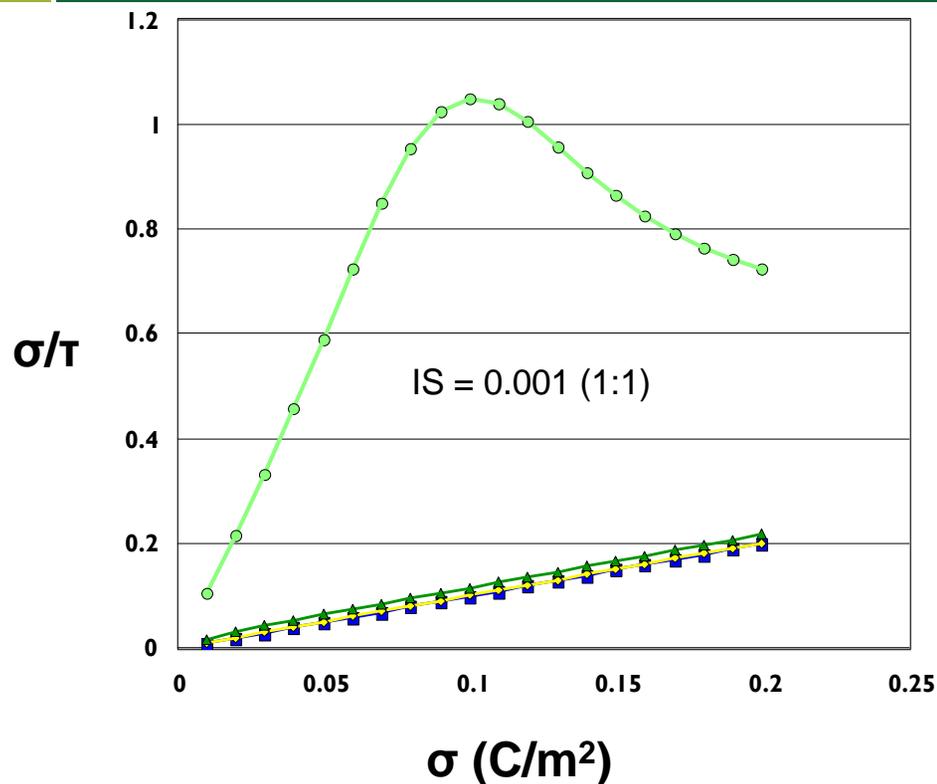
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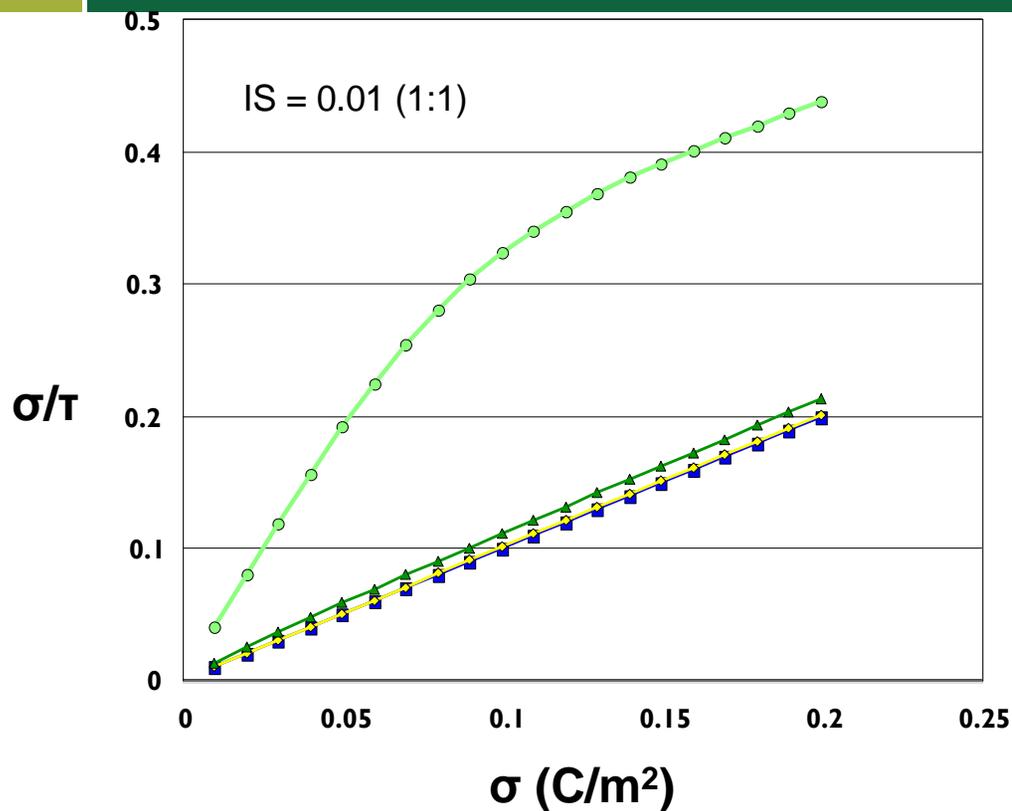


Michelis-Menten plots: $r_{pcl} = 1 \mu\text{m}$ (blue boxes), $0.1 \mu\text{m}$ (yellow diamonds), $0.01 \mu\text{m}$ (green triangles); and $0.001 \mu\text{m}$ (aquamarine circles). Slope = $1/\tau_{max}$.

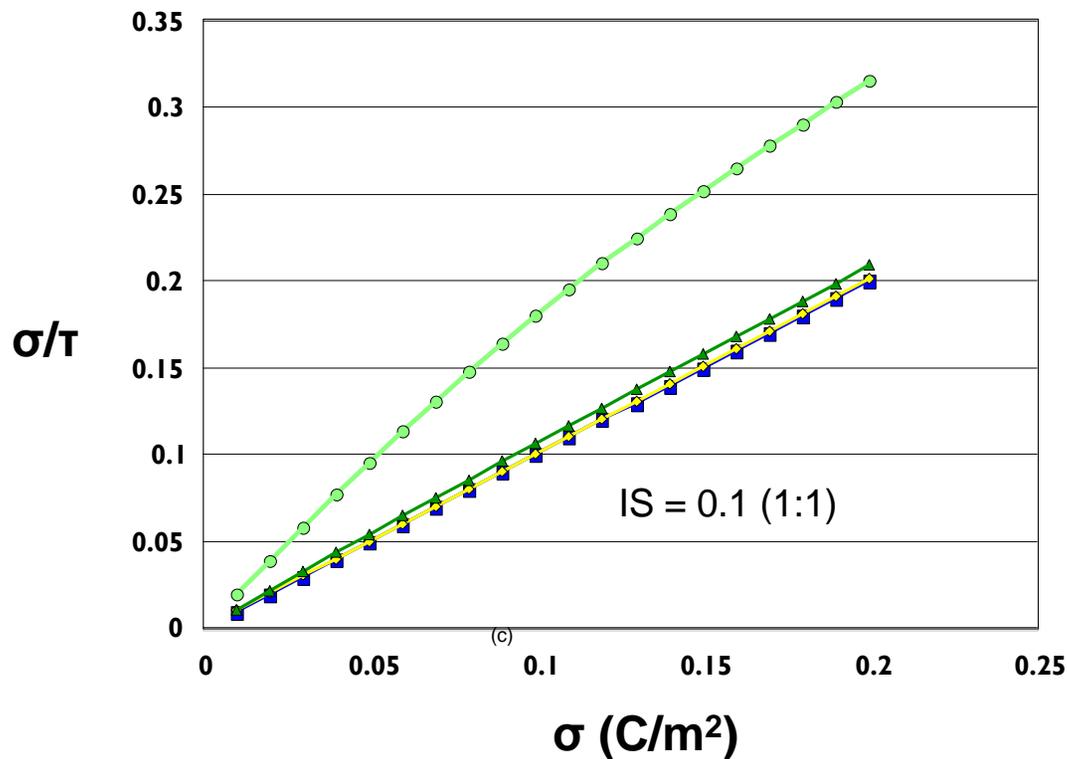


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Michelis-Menten plots: $r_{pci} = 1 \mu\text{m}$ (blue boxes), $0.1 \mu\text{m}$ (yellow diamonds), $0.01 \mu\text{m}$ (green triangles); and $0.001 \mu\text{m}$ (aquamarine circles). Slope = $1/\tau_{max}$.

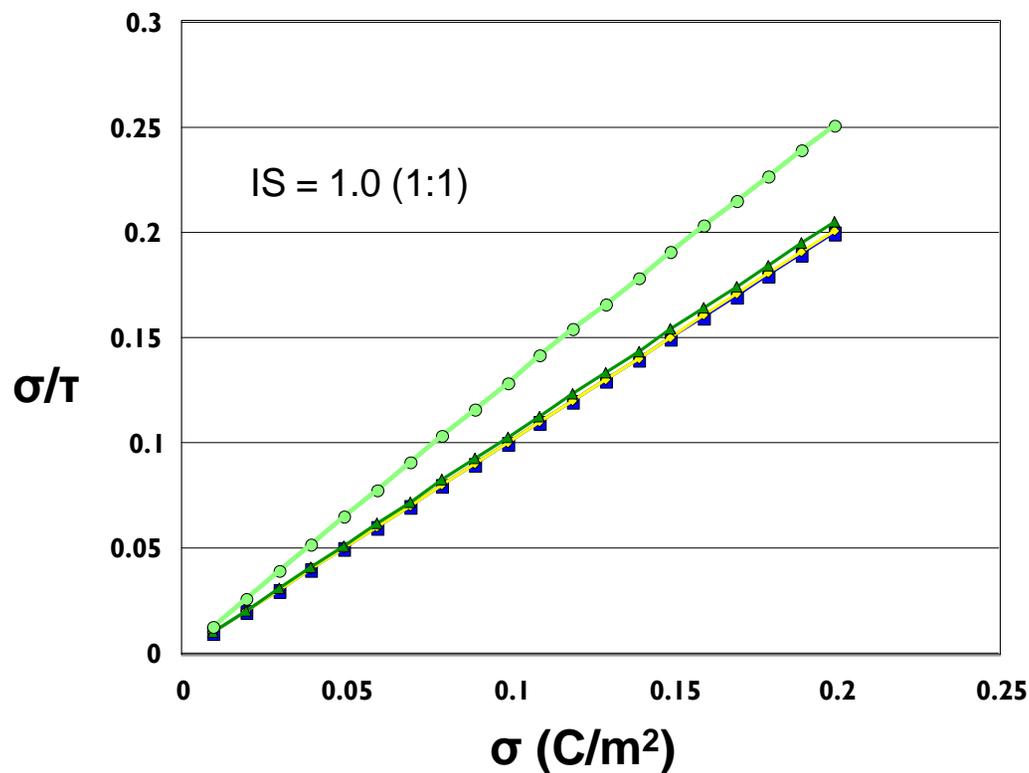


Michelis-Menten plots: $r_{pci} = 1 \mu\text{m}$ (blue boxes), $0.1 \mu\text{m}$ (yellow diamonds), $0.01 \mu\text{m}$ (green triangles); and $0.001 \mu\text{m}$ (aquamarine circles). Slope = $1/\tau_{max}$.



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$\log_{10}(IS)$ ----- Estimates of τ_{max} -----

	$r_{pcl} =$ 1 μm	$r_{pcl} =$ 0.1 μm	$r_{pcl} =$ 0.01 μm	$r_{pcl} =$ 0.001 μm
-3	0.9996	0.9955	0.9640	~0.4
-2	0.9996	0.9956	0.9569	~0.5
-1	0.9996	0.9963	0.9629	~0.6
0	0.9998	0.9977	0.9774	0.7990

$\Delta G_{charging} < 0.1 e\psi$

$\Delta G_{charging} > 0.1 e\psi$



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Conclusions

- *In both the present approach and the DHLL, the principal excess free energy term is an ionic-strength-dependent charging energy term.*
- *The historical mass action expression is supported in systems where $r_{pcl} \geq 0.1 \text{ um}$ and $IS \geq 0.001 \text{ M (1:1)}$.*
- *At high σ , the historical expression is supported when $r_{pcl} \geq 0.01 \text{ um}$ and $IS \geq 0.001 \text{ M (1:1)}$.*
- *Charging energies are predicted to be significant in all situations where $r_{pcl} \approx 0.001 \text{ um}$ (i.e., charged bound sites are predicted to display non-ideal behavior).*

*How significant will be the **red** term with smaller particles in intermediate charge density situations?*

$$\Psi_d = \frac{(1-\tau)Q_p}{4\pi\epsilon r_{pcl}} - \frac{(1-\tau)Q_p}{4\pi\epsilon(r_{pcl} + 1/\kappa)}$$

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$$\Psi_d = \frac{(1-\tau)Q_p}{4\pi\epsilon r_{pcl}} - \frac{(1-\tau)Q_p}{4\pi\epsilon(r_{pcl} + 1/\kappa)}$$

Specifically:

$$\ln(\gamma) = \frac{+(1-\tau)(\pm 2eQ_p + e^2)}{(r_{pcl} + 1/\kappa)8\pi\epsilon kT} - \frac{(1-\tau)(\pm 2eQ_p + e^2)}{r_{pcl}8\pi\epsilon kT}$$

??????